gnals on heating at ~5 deg ide with Yoder's data, within ion. Other high-pressure work wersion has been noted by \$\mathcal{F}\$6].

-- ruartz-coesite triple point erous investigations of the ration [MacDonald, 1956] , 1959; Boyd and England 15-3; Takahashi, 1963; Kita-1964; Bell et al., 1965; . the following data are sesition at high temperatures: position is the 1400°C, 37.5 estained with Pt versus Pt mocouples by Boyd et al. relative set of data is that of [1960] as revised [Boud thing these data with a slight te present results (Figure 1). =-quartz-coesite triple point and 37 kb.

f the hypothesis of a first-Investigations of the lowsizion are very numerous: unclear that definitive and we yet been obtained because and rapid variation of therleters near the inversion. For et al. [1953] report a fivehe coefficient of thermal exmerval of less than 2°, slightly ition temperature. Detailed intercomparison of heat cagransion, and elastic moduli thermodynamic consequences second paper of this series iten, in preparation). It apntortant that these measure-Eder isothermal conditions, on hest purity.

tors have considered the highion to be a first-order transines for the discontinuous in- ΔV , upon heating through the on include (in cubic centirula weight (fw)): ~ 0.195 evaluating earlier measure- $Mz[umdar\ et\ al.,\ 1964]$; 0.154 cl., 1966]. Some estimates for ge. ΔS , of the transition in- τ degree per formula weight): ~0.74 [Sosman, 1927, evaluating earlier measurements]; ~0.43 [Majumdar et al., 1964]. Berger et al. [1965] used thermogrammatic rechniques on a variety of samples and obtained values ranging from ~0.5 to ~1.0 J \log^{-1} fw⁻³ (Kelley [1960] tabulates a finite ΔS for the inversion, but this is a manner of presenting the data and does not mean to imply that the transition is considered as first order (K. K. Kelley, personal communication).)

The Clausius-Clapeyron equation, $dT/dp = \Delta V/\Delta S$, constrains estimates of ΔV and ΔS to be consistent with the initial slope of the phase boundary, $dT/dp = 26 \pm 1 \text{ deg kb}^{-1}$. For any erst-order transition, there is a further constraint in that the initial curvature

$$\frac{f^{2}T}{d\rho^{2}} = \frac{1}{\Delta S} \left\{ \left(\frac{\partial \Delta V}{\partial \rho} \right)_{T} + 2 \left(\frac{dT}{d\rho} \right) \left(\frac{\partial \Delta V}{\partial T} \right)_{p} - \left(\frac{dT}{d\rho} \right)^{2} \frac{\Delta C_{p}}{T} \right\}$$
(1)

For the present results, $-d^2T/dp^2 \lesssim 0.4$ deg kb⁻². No useful application of the Clausius-Clapeyron equation to the high-low quartz transition can be made with ΔS and ΔV values that vary over a factor of 2. Initial slopes calculated from the quantities tabulated span the experimentally determined value. Application of (1) or even Bridgman's [1931] empirical correlations,

$$-\left(\frac{\partial \Delta V}{\partial p}\right)_{T} \ge \left(\frac{dT}{dp}\right) \left(\frac{\partial \Delta V}{\partial T}\right)_{p} \ge \left(\frac{dT}{dp}\right)^{2} \left(\frac{\Delta C_{p}}{T}\right) \tag{2}$$

s made most uncertain because of the rapid variations in $(\partial V/\partial p)_T$, $(\partial V/\partial T)_p$, and C_p as the transition is approached. Various bounds may be obtained via (2), using published estimates for discontinuities in C_p/T [e.g., Berger t al., 1965] or $(\partial V/\partial T)_p$ [e.g., Sosman, 1927; Rosenholtz and Smith, 1941], but the worth of any of these values for the possible discontinuities is very much in doubt.

Perhaps one of the most experimentally accessible constraints on the possible first-order transition is that involving trajectories of phase boundaries at the high-quartz-low-quartz-coesite triple point. Using 2.93 g cm^{-a} as the probable density of coesite at room temperature and tero pressure [Frondel, 1962, pp. 310 and 314].

 ΔV (low quartz \rightarrow coesite) ~ 2.2 cm³ fw⁻¹ and a similar value may be assumed near the triple point. The 'averaged dT/dp,' ('quartz-coesite) is ~ 9 deg kb⁻¹ [Boyd and England, 1960, et seq.]. To a good approximation, the difference in slope between the high-quartz-coesite and low-quartz-coesite phase boundaries, dT/dp (high-quartz-coesite) -dT/dp (low-quartz-coesite) $\equiv \Delta dT/dp$, is

$$\Delta \frac{dT}{dp} \approx \frac{\Delta V(\text{low} \to \text{high quartz})}{\Delta V(\text{low quartz} \to \text{coesite})} \cdot \frac{dT}{dp} \text{ ('quartz'-coesite)}$$
(3)

From (3) and an average value (see above) for ΔV (low \rightarrow high quartz), $-\Delta \ dT/dp \sim 0.6$ deg kb⁻¹. The best high-pressure data presently available [Boyd and England, 1960] are inadequate, however, to detect a kink of such magnitude in the quartz-coesite phase boundary.

In the authors' opinion, conclusive proof for first-order behavior of the high-low quartz transition is lacking. In any case, the most spectacular variations in the thermophysical properties near the inversion indicate a lambda transition, the thermodynamics of which is discussed in a forthcoming paper (Klement and Cohen, in preparation).

Geological implications. For the several temperature distributions within the crust and upper mantle calculated by Clark and Ringwood [1964], all the proposed temperaturedepth curves lie within the field of low quartz; nevertheless, regions of significantly elevated temperatures, as well as magmas, may be well within the stability field of high quartz. A number of equilibrium reactions involving quartz intersect the low-high quartz inversion curve. If definite criteria were available to distinguish quartz that has passed through the inversion from quartz that has not, then comparison of reaction curves with the data for quartz may give additional constraints on the paragenesis of the assemblage being considered. Frondel [1962, p. 119] has summarized the problem of inversion criteria as follows: 'The question arises whether a given natural quartz crystal or anhedron originally crystallized as low-quartz or as high-quartz. . . . In general the criteria as afforded by natural material are